## Defect Model of a Tetragonal Sm<sup>3+</sup> Center Found from EPR Measurements in CaF<sub>2</sub> and SrF<sub>2</sub> Crystals

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The EPR parameters (**g** factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$ ) of a tetragonal (C<sub>4v</sub>) Sm<sup>3+</sup> center in CaF<sub>2</sub> and SrF<sub>2</sub> crystals are calculated by considering the crystal-field J-mixing among the ground  $^6{\rm H}_{5/2}$ , the first excited  $^6{\rm H}_{7/2}$  and second excited  $^6{\rm H}_{9/2}$  state multiplets. In the calculations the free-ion and crystal-field parameters of the tetragonal Sm<sup>3+</sup>-F<sup>-</sup> center obtained from polarized laser-selective excitation spectroscopy are used. The calculated results suggest that the tetragonal Sm<sup>3+</sup>-F<sup>-</sup> center is the Sm<sup>3+</sup> center found by later EPR measurements. The **g** factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$  of this EPR center are satisfactorily explained.

*Key words:* Defect Model; Electron Paramagnetic Resonance; Crystal-field Theory; Sm<sup>3+</sup>; CaF<sub>2</sub>; SrF<sub>2</sub>.

## 1. Introduction

The doping of alkaline earth fluorides CaF2 and SrF<sub>2</sub> with rare earth ions usually results in the replacement of divalent alkaline ions by trivalent rare earth (Re<sup>3+</sup>) ions. The required charge compensation can occur in many ways, leading to sites with cubic, trigonal, tetragonal or rhombic symmetry [1, 2]. There may be different impurity centers with even a consistent axial (C<sub>4v</sub> or C<sub>3v</sub>) symmetry in Re<sup>3+</sup>-doped CaF2 and SrF2 crystals. For example, early EPR studies [1,3] found a tetragonal  $(C_{4v})$  Sm<sup>3+</sup> center with  $g_{\parallel} \approx 0.907(10)$  and  $g_{\perp} \approx 0.544(10)$  in CaF<sub>2</sub>: Sm<sup>3+</sup>. Weber and Bierig [1] suggested that this center probably arises from an interstitial F- charge compensator located at the center of one of the nearest empty cubes of the lattice (so we name it  $Sm^{3+}$ - $F^-$  ( $C_{4v}$ ) center). On later studies [4-7] of EPR in CaF<sub>2</sub>: Sm<sup>3+</sup> and SrF<sub>2</sub>:Sm<sup>3+</sup> a consistent C<sub>4v</sub> symmetry center was found that was spectroscopically distinct (e.g.,  $g_{\parallel} \approx$  $0 \pm 0.06$ ,  $g_{\perp} \approx 0.823 \pm 0.003$  in CaF<sub>2</sub>:Sm<sup>3+</sup>) and thus differed in the charge compensation configuration from the earlier studies. However, this  $Sm^{3+}(C_{4v})$  center was also suggested as  $Sm^{3+}$ - $F^-$  ( $C_{4v}$ ) center [7, 8]. It

is interest to determine which EPR center is the Sm<sup>3+</sup>-F<sup>-</sup> (C<sub>4v</sub>) center. Polarized laser-selective excitation is a well-established method for identifying lines of multicenter spectra [9]. Polarized laser-selective excitation and fluorescence spectroscopy of Sm<sup>3+</sup>-doped CaF<sub>2</sub> and SrF2 crystals was recently performed by Wells and Reeves [10]. They [10] found that the dominant center present in both host crystals is the Sm<sup>3+</sup>-F<sup>-</sup>  $(C_{4v})$  center and the optical spectral parameters (i. e., the free-ion and crystal-field parameters) of this center in both crystals were obtained (see Table 1). Although they pointed out that there are two different  $\mathrm{Sm}^{3+}$  ( $\mathrm{C}_{4\mathrm{v}}$ ) centers found by the EPR measurements, they did not suggest which EPR Sm<sup>3+</sup> (C<sub>4v</sub>) center is the  $Sm^{3+}$ - $F^-$  ( $C_{4v}$ ) center. Since the EPR parameters (g factors and hyperfine structure constants A) of a 4f<sup>n</sup> ion in low symmetry are sensitive to the optical spectral parameters, in this paper we have calculated the EPR g factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$  for a tetragonal Sm $^{3+}$  center in CaF2 and SrF2 crystals by using the above optical spectral parameters. The results (which are related to the assignment of the EPR C<sub>4v</sub> center) are discussed.

Table 1. Free-ion and crystal-field parameters (in cm<sup>-1</sup>) for the Sm<sup>3+</sup>-F<sup>-</sup> (C<sub>4v</sub>) centers in CaF<sub>2</sub> and SrF<sub>2</sub> crystals [10].

Parameter	CaF <sub>2</sub> :Sm <sup>3+</sup>	SrF <sub>2</sub> :Sm <sup>3+</sup>	
$F^2$ $F^4$	78824	79062	
$F^4$	56842	56870	
$F^6$	39972	40077	
$\alpha$	20.6	20.6	
β	-724	-724	
r	1700	1700	
$\zeta_{ m 4f}$	1166	1168	
$B_A^2$	746	472	
$B_A^4$	590	545	
$B_A^6$	617	489	
$egin{array}{c} egin{array}{c} egin{array}{c} egin{array}{c} B_A^4 \ B_A^6 \ B_C^2 \ B_C^4 \ \end{array}$	-1227	-1204	
$B_C^4$	670	562	

## 2. Calculations

The ground state of a free Sm<sup>3+</sup> (4f<sup>5</sup>) ion is  ${}^{6}H_{5/2}$ , which is split into three Kramers doublets in a tetragonal crystal field. The lowest (or ground) doublet  $\Gamma \gamma$ may be  $\Gamma_6$  or  $\Gamma_7$  depending upon the crystal field parameters. Since the g factors calculated by considering the mixing in only the ground state multiplet <sup>6</sup>H<sub>5/2</sub> or, further, the crystal-field J-mixing of the first excited state multiplet <sup>6</sup>H<sub>7/2</sub> into the ground state multiplet  ${}^{6}\mathrm{H}_{5/2}$  can not agree with the observed values [11], we consider the J-mixing among the ground  ${}^{6}H_{5/2}$ , the first excited  ${}^{6}H_{7/2}$  and second excited  ${}^{6}H_{9/2}$  state multiplets via crystal-field interaction here. Thus a  $24 \times 24$ energy matrix is established. Substituting the free-ion and crystal-field parameters obtained from the optical spectra for the  $Sm^{3+}$ - $F^-$  ( $C_{4v}$ ) center in  $CaF_2$  and  $SrF_2$ crystals (see Table 1) into the matrix and diagonalizing it, we can obtain that the wave function of the ground doublet of CaF<sub>2</sub>:Sm<sup>3+</sup> is

$$\begin{split} |\varGamma_6\rangle \approx & \pm 0.9865 |5/2, \pm 1/2\rangle - 0.0413 |7/2, \pm 1/2\rangle \\ & - 0.0867 |7/2, \mp 7/2\rangle \pm 0.0037 |9/2, \pm 1/2\rangle \quad (1) \\ & \pm 0.1251 |9/2, \mp 7/2\rangle \pm 0.0052 |9/2, \pm 9/2\rangle, \\ \text{and that for SrF}_2: Sm^{3+} \text{ it is} \\ |\varGamma_6\rangle \approx & \pm 0.9890 |5/2, \pm 1/2\rangle - 0.0326 |7/2, \pm 1/2\rangle \\ & - 0.0890 |7/2, \mp 7/2\rangle \pm 0.0022 |9/2, \pm 1/2\rangle \quad (2) \\ & \pm 0.1130 |9/2, \mp 7/2\rangle \pm 0.0103 |9/2, \pm 9/2\rangle. \end{split}$$

From the Zeeman interaction  $H_z$  (=  $g_J \mu_\beta \mathbf{H} \cdot \mathbf{J}$ , with the original meanings [12, 13]) and hyperfine interaction  $H_{\rm hf}$  (=  $PN_J\hat{N}$ , where P is the dipolar hyperfine

structure constant and  $N_J$  the diagonal matrix element for the  $^{2S+1}L_J$  state [12]), we have the perturbation formulas of the EPR parameters for  $4f^n$  ions as follows:

$$g_{\parallel} = 2g_{J}\langle \Gamma \gamma | \hat{J}_{Z} | \Gamma \gamma \rangle, \ g_{\perp} = g_{J}\langle \Gamma \gamma | \hat{J}_{+} | \Gamma \gamma' \rangle,$$
$$A_{\parallel} = 2PN_{J}\langle \Gamma \gamma | \hat{N}_{Z} | \Gamma \gamma \rangle, \ A_{\perp} = PN_{J}\langle \Gamma \gamma | \hat{N}_{+} | \Gamma \gamma' \rangle. \tag{3}$$

Considering the covalence of the Sm3+-F- bond in both crystals, the orbital angular momentum  $\hat{L}$  in the above formulas should be multiplied by an orbit reduction factor k. This factor, which is slightly smaller than 1, depends on the covalence of the metal-ligand bond. The shorter the metal-ligand distance R, the stronger the covalence of this bond and so the smaller the factor k. This point can be confirmed by the following facts: (i) The covalence of the host crystal CaF<sub>2</sub> is slightly stronger than that of the host crystal SrF<sub>2</sub> [14]. (ii) The effect of pressure on the freeion parameters (Coulomb repulsion  $F^K$  and spin-orbit coupling coefficient  $\zeta_{4f}$ ) of the Re<sup>3+</sup> ions in crystal [15, 16] suggests that these parameters decrease with decreasing metal-ligand distance and hence with increasing covalence of the bond. For CaF<sub>2</sub>:Sm<sup>3+</sup>, the above free-ion parameters are slightly smaller than those of SrF<sub>2</sub>:Sm<sup>3+</sup> (see Table 1), and so its covalence is stronger. Thus we can reasonably assume for CaF<sub>2</sub>:Sm<sup>3+</sup>,  $k \approx 0.976$  and for SrF<sub>2</sub>:Sm<sup>3+</sup>,  $k \approx 0.980$ . Applying (1) and (2), the factors k and the free-ion values of P ( $^{147}$ Sm)  $\approx -51.7(6) \times 10^{-4}$ cm<sup>-1</sup> and P  $(^{149}\text{Sm}) \approx -41.8(6) \times 10^{-4}\text{cm}^{-1}$  [12] to (3), the EPR parameters  $g_{\parallel}$ ,  $g_{\perp}$ ,  $A_{\parallel}$  and  $A_{\perp}$  for the tetragonal Sm<sup>3+</sup> centers in both crystals are calculated. The results are compared with the observed values in Table 2.

## 3. Conclusion and Discussion

From Table 2, it can be seen that the calculated EPR parameters using the free- ion and crystal-field parameters obtained from the optical spectra of Sm³+- F- (C<sub>4v</sub>) centers in CaF<sub>2</sub> and SrF<sub>2</sub> crystals are consistent with the observed values given in [4–7]. So, the tetragonal EPR Sm³+ center in CaF<sub>2</sub> and SrF<sub>2</sub> found in [4–7] rather than that in [1,3] is the Sm³+-F- (C<sub>4v</sub>) center. The assignment can not be transformed by changing the orbit reduction factor k because the factor k affects slightly the calculated average value of  $\bar{g}[=(g_{\parallel}+2g_{\perp})/3]$  and  $\bar{A}[=(A_{\parallel}+2A_{\perp})/3]$ , but it can not alter the sign of the anisotropy of the  ${\bf g}$  factor and constant A (characterized by  $\Delta g = g_{\parallel} - g_{\perp}$  and  $\Delta A = A_{\parallel} - A_{\perp}$ ). So, the above assignment is reasonable

		$g_{\parallel}$	$g_{\perp}$	$A_{\parallel}$ (147Sm)	$A_{\perp}(^{147}\mathrm{Sm})$	$A_{\parallel}$ (149Sm)	$A_{\perp}(^{149}\mathrm{Sm})$
CaF <sub>2</sub> :Sm <sup>3+</sup>			$0.822 \\ 0.823 \pm 0.003 \\ 0.544 \pm 0.010$		$221 \pm 5$ $230 \pm 5$	6.1 0 ± 10	$180 \pm 4$ $190 \pm 5$
SrF <sub>2</sub> :Sm <sup>3+</sup>	Cal. Expt. [4-7]	0.027 < 0.1	$0.824$ $0.823 \pm 0.003$	7.3 0±10	$225 \pm 5$ $230 \pm 4$	$6.2 \\ 0 \pm 10$	$183 \pm 4$ $190 \pm 4$

Table 2. The EPR  $\mathbf{g}$  factors and hyperfine structure constants  $A_i$  (in units of  $10^{-4} \text{cm}^{-1}$ ) for the tetragonal Sm<sup>3+</sup> center in CaF<sub>2</sub> and SrF<sub>2</sub> crystals.

and the **g** factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$  (for isotopes <sup>147</sup>Sm<sup>3+</sup> and <sup>149</sup>Sm<sup>3+</sup>) observed in [4–7] for Sm<sup>3+</sup>-doped CaF<sub>2</sub> and SrF<sub>2</sub> can be

satisfactorily explained. The defect model of the tetragonal EPR  $\rm Sm^{3+}$  center in  $\rm CaF_2$  found in [1, 3] remains to be further studied.

- [1] M. J. Weber and R. W. Bierig, Phys. Rev. 134A, 1492 (1964).
- [2] T. S. Chang and M. I. Cohen, J. Chem. Phys. 64, 5255 (1976).
- [3] W. Low, Phys. Rev. 134A, 1479 (1964).
- [4] A. A. Antinpin, I. N. Kurkin, L. D. Livanova, L. Z. Potvorova, and L. Ya. Shekun, Sov. Phys. Solid State 7, 1271 (1965).
- [5] A. A. Antinpin, I. N. Kurkin, L. D. Livanova, L. Z. Potvorova, and L. Ya. Shekun, Sov. Phys. Tech. Phys. 11, 821 (1966).
- [6] H. N. Evans and S. D. Mclaughlan, Phys. Lett. 23, 638 (1966).
- [7] R. C. Newman and R. J. Woodward, J. Phys. C 7, L433 (1974).
- [8] I. J. Ashburner, R. C. Newman, and S. D. Mclaughlan, Phys. Lett. 27A, 212 (1968).

- [9] J. C. Wright and K. M. Cirillo-Penn, Radiat. Eff. Def. Solids, 119 – 121, 231 (1991).
- [10] J.-P. R. Wells and R. J. Reeves, Phys. Rev. B61, 13593 (2000).
- [11] M. Yamaga, M. Honda, J.-P.R. Wells, T. P.J. Han, and H. G. Gallagher, J. Phys.:Condens. Matter 12, 8727 (2000).
- [12] A. Abragam and B. Bleanely, Electron Paramagnetic Resonance of Transition-Ions, Oxford University press, London 1970.
- [13] I. A. Sorin and M. V. Vlasova, Electron Spin Resonance of Paramagnetic Crystals (Translated from Russian by P. Gluck), Plenum Press, New York 1973.
- [14] B. F. Levine, J. Chem. Phys. 59, 1463 (1973).
- [15] C. Bungenstock, Th. Troster, and W.B. Holzapfel, Phys. Rev. B62, 7945 (2000).
- [16] Y. R. Shen and W. B. Holzapfel, Phys. Rev. **B52**, 12618 (1995).